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Title: Image Station Use Examples from DARHT

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Image Station Use Examples from DARHT

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Los Alamos National Laboratory

ASD-Scorpius Downstream Transport Meeting
January 29, 2021

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Outline of Material

- The beam envelope equation for tuning
- Optical Transition Radiation (OTR)
- Solenoid scan method for beam identification
- Pepper-pot method for beam identification
- Beam-target studies via gated imaging

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A word about sources...

**“Immature artists imitate.
Mature artists steal.”**

-Various attributions (Eliot, Picasso, Trilling, others)*

- J-6 scientists are the local experts on these techniques (DC Moir, C Ekdahl, M Schulze)
- This talk is mostly a summary of what they've been doing for years (and I've just been learning)

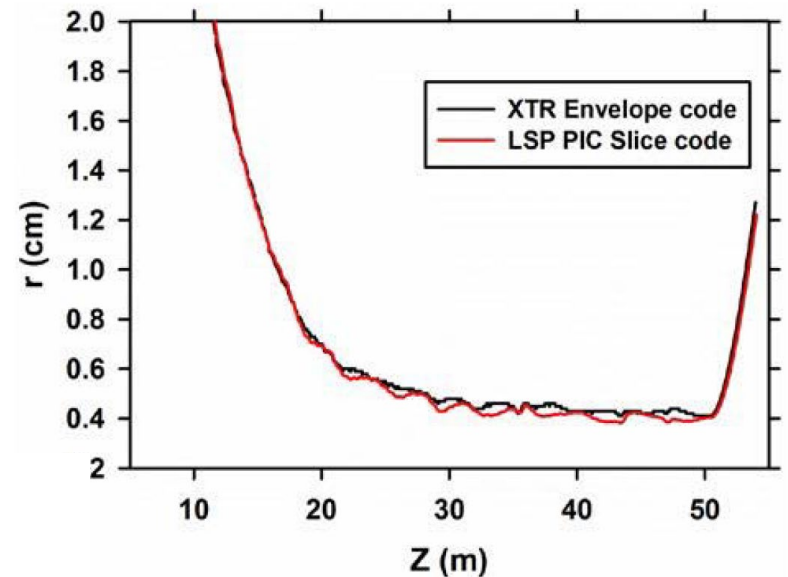
*Quote Investigator <https://quoteinvestigator.com/2013/03/06/artists-steal/>

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Validated theory enables design

- Envelope equation describes beam radius evolution
- Various derivations either from beam moments¹ or paraxial equations²
- Use for LIA analysis developed to high degree by P. Allison in xtr³
- Checked against various PIC codes over the years⁴

$$r_m'' + \underbrace{\frac{\gamma' r_m'}{\beta^2 \gamma}}_{\text{Accel. damping}} + \underbrace{\frac{\gamma'' r_m}{2\beta^2 \gamma}}_{\text{Gap accel.}} + \underbrace{\left(\frac{qB}{2mc\beta\gamma}\right)^2 r_m}_{\text{Sol. focus}} - \underbrace{\left(\frac{p_\theta}{mc\beta\gamma}\right)^2 \frac{1}{r_m^3}}_{\text{Angular momentum}} - \underbrace{\frac{\epsilon_n^2}{\beta^2 \gamma^2 r_m^3}}_{\text{Emittance}} - \underbrace{\frac{K}{r_m}}_{\text{Self-fields}} = 0.$$



Ekdahl, et al., IEEE TPS (2017)

¹Lee and Cooper, Part. Accel. 7 (1976) 83.

²cf Reiser "Theory and Design..." 2008. and Humphries "Charged Particle Beams" 2002.

³Allison, LA-UR-01-6585 (2001)

⁴Ekdahl, et al., IEEE Trans. Plasma Sci. 45 (2017) 2962.

Radius, convergence, and emittance are key parameters in the envelope equation

$$r_m'' + \frac{\gamma' r_m'}{\beta^2 \gamma} + \frac{\gamma'' r_m}{2\beta^2 \gamma} + \left(\frac{qB}{2mc\beta\gamma} \right)^2 r_m - \left(\frac{p_\theta}{mc\beta\gamma} \right)^2 \frac{1}{r_m^3} - \frac{\epsilon_n^2}{\beta^2 \gamma^2 r_m^3} - \frac{K}{r_m} = 0.$$

- Equation is second order and requires initial conditions R_0 , dR_0/dz and ϵ_n to solve
- Constant emittance is assumed, though models exist to describe growth
- B-fields and acceleration gaps are considered knowns with energy & current
- Equation is symmetric in z: can be solved “backwards”
- Uncertainties propagate or may only weakly determine beam properties far away from measurement point

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Outline of Material

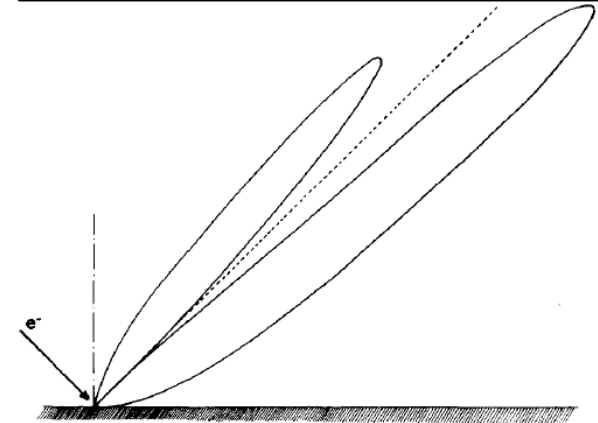
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- Beam-target studies via gated imaging

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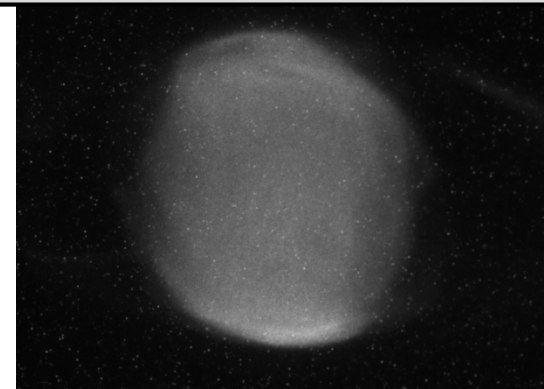
Optical transition radiation is produced where electrons transit material boundaries

- OTR can be considered a special form of Cerenkov radiation^{1,2}
- Radiation emitted when const. velocity particle traverses materials with different dielectric constants
- Simple and complex uses exist in literature
 - “Near-field” images source location and intensity
 - “Far-field” can extract convergence and emittance

Radiation Pattern, $\gamma=10$ ¹



DARHT-2 ISC Image (2019)



¹Wartski, et al., J. Appl. Phys. 46 (1975) 3644.

²Fiorito and Rule, Conf. Proc. Beam Instrum. Workshop (1994) 21.

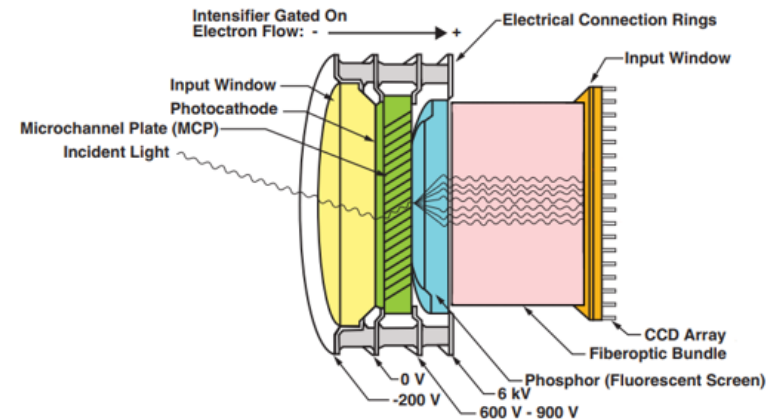
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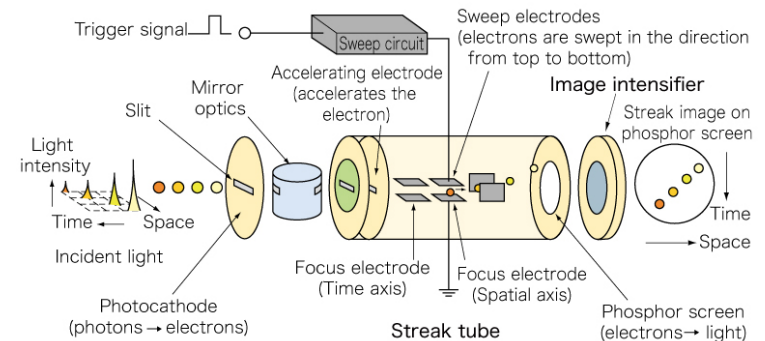
Time-resolved or gated imaging is critical to these diagnostics

- Beam dynamics and target phenomena evolve on \sim ns time scale
- Intensified imagers amplify photon-starved signals
- Gated cameras retain 2D information for a single time window
- Streak systems can provide continuous, time-resolved 1D “slice”

Intensified CCD Operating Principle



Streak Camera Operating Principle

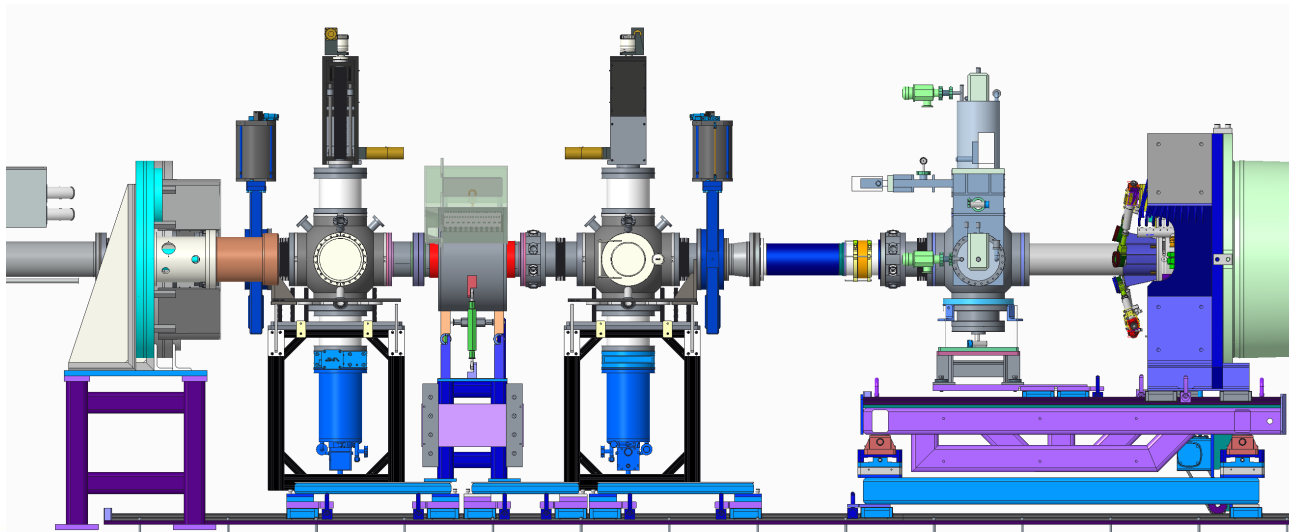


¹<https://www.princetoninstruments.com/learn/camera-fundamentals/iccd-and-emiccd-basics>

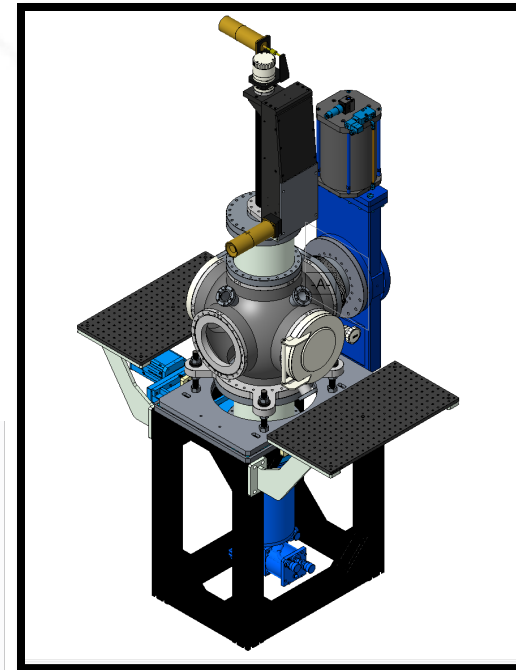
²<https://www.hamamatsu.com/us/en/product/photometry-systems/streak-camera/operating-principle/index.html>

OTR alone at two locations can provide **R** and **R'**

- Direct measure of R with known z between locations
- Single camera with mirrors or
- Simultaneous measurements (with caveats for foil focusing + scattering)



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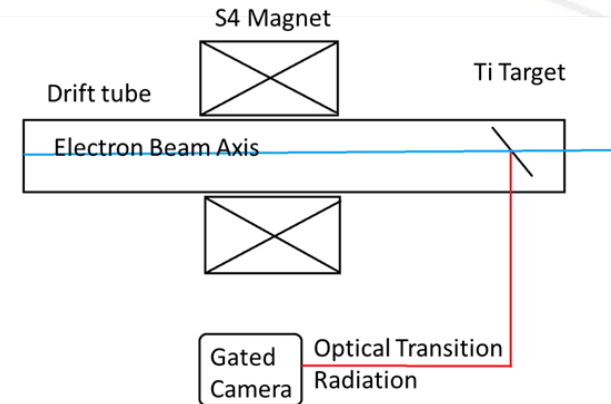
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OTR + solenoid scans can extract R , R' , and ϵ_n

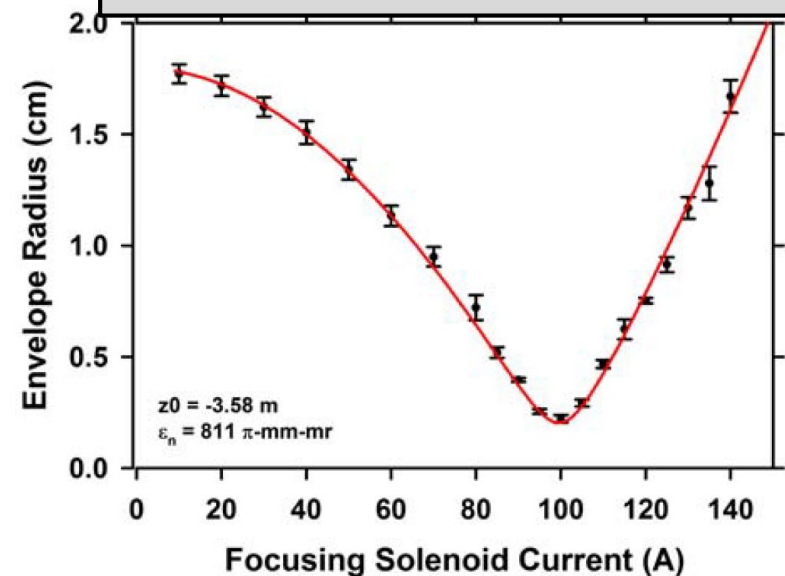
$$r_m'' + \left(\frac{qB}{2mc\beta\gamma} \right)^2 r_m - \frac{\epsilon_n^2}{\beta^2\gamma^2 r_m^3} - \frac{K}{r_m} = 0$$

- Envelope through magnetic field can be solved for entrance conditions (R , R' , ϵ_n)
- Scan in magnetic field changes waist location: sum of data provides unique solution to three variables
- Zeroth order consideration: don't destroy the target and avoid beam-target interaction effects!
 - Long focal length increases spot size
 - Short pulses and gating early in the pulse reduces light-ion effects

Typical Configuration



DARHT-2 Solenoid Scan Result



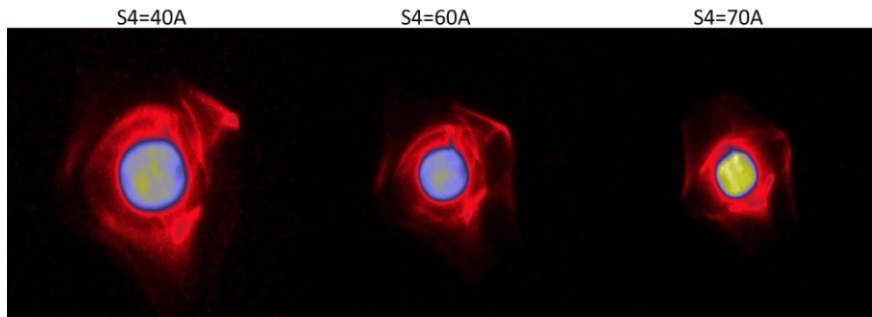
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Additional corrections and interpretations

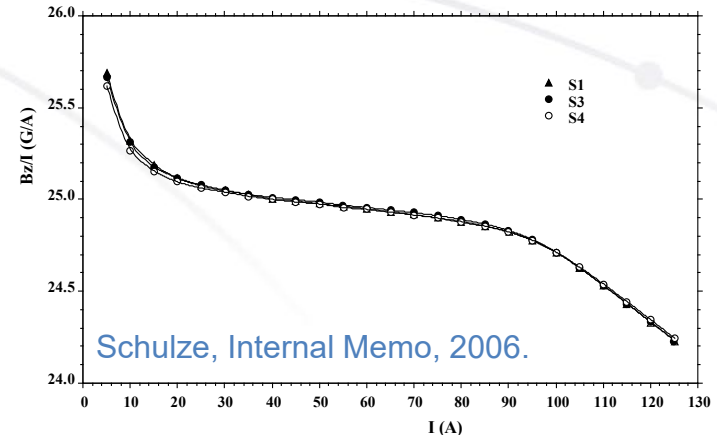
- Magnetic field is assumed “known” but non-linear excitation affects analysis
- High-contrast imaging can reveal beam halo and indicate larger ϵ_n (Moir and Allison, LA-UR-21-21386)
- Spherical aberrations can modify results as well (Schulze LA-UR-20-24545)

Beam Halo

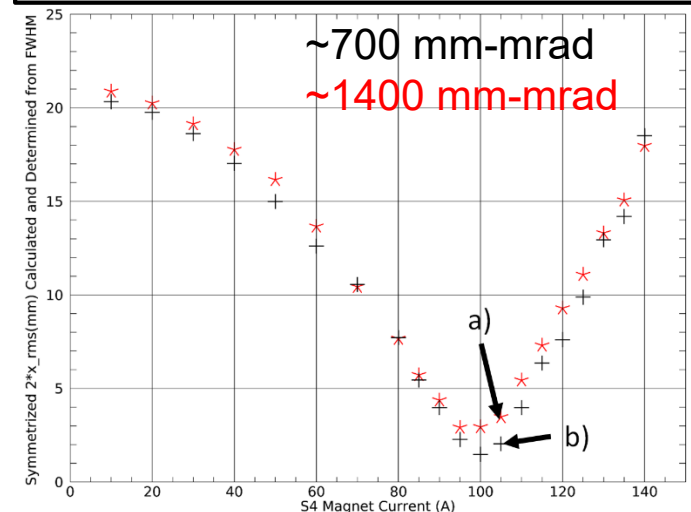


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Magnet Excitation Fall-off



Solenoid Scan w/ and w/o Halo



Moir and Allison, LA-UR-21-21386

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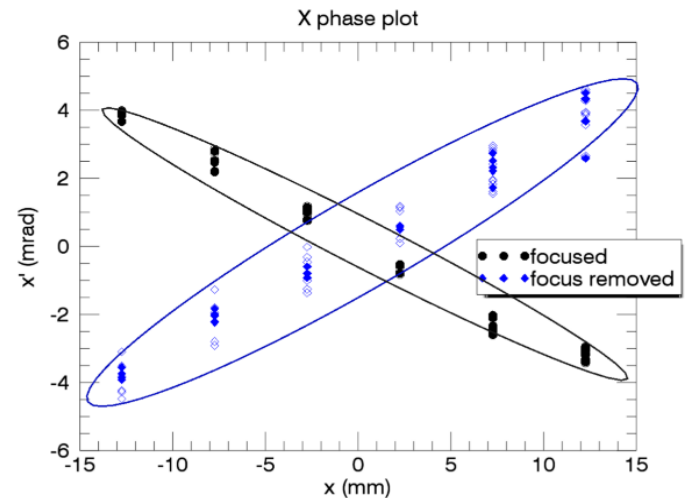
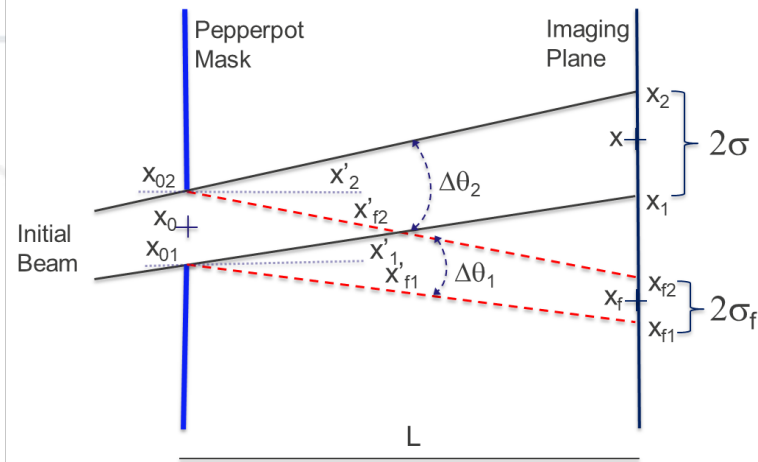
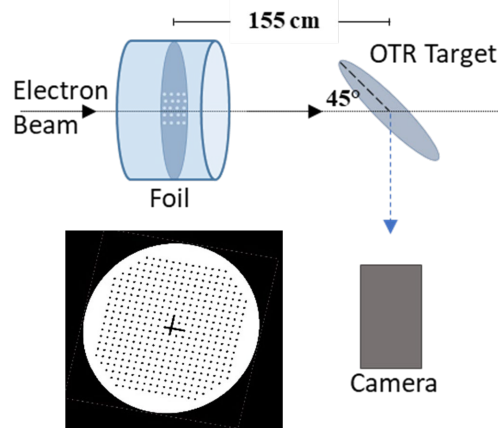
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Pepper-pot diagnostic method can extract ϵ_n of the beam

- Thin pepper-pot creates contrast by scattering portions of the beam
- Beamlet size relative to mask gives emittance
- In principle:* eliminates need for dual-imaging and reduces beam-target interactions



in π -mm-mrad	ϵ_x	ϵ_y
Focused	147	306
Corrected	226	415

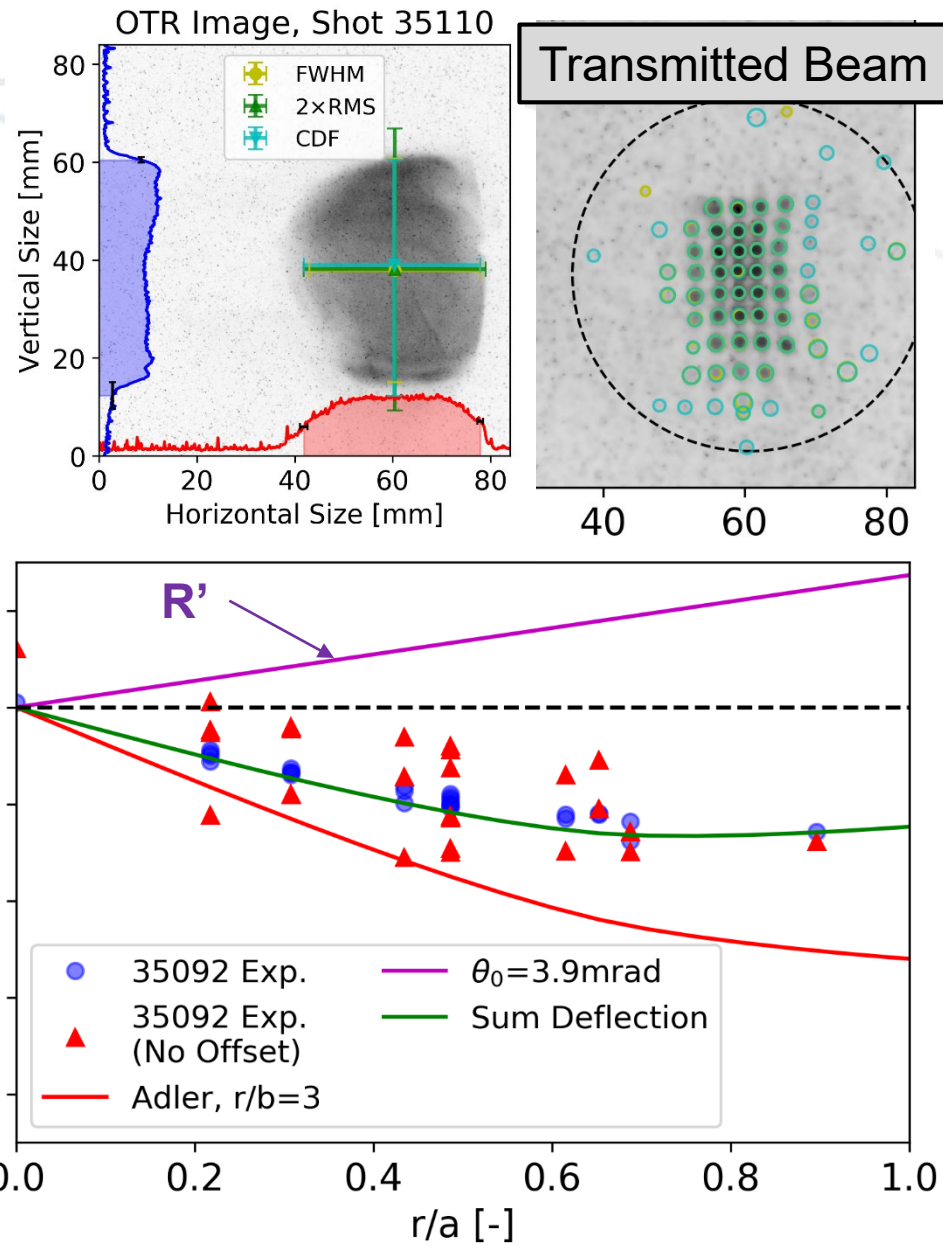
Schultz, et al., Proc. NAPAC 2019

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Single pepper-pot can also provide **R**, and **R'** measurements

- **R** can be determined from counting mask dots (low resolution)
- **R'** from beam distortion and foil-focusing
- ϵ_n from spot-size analysis
- *Not really this simple: analysis still in development!*



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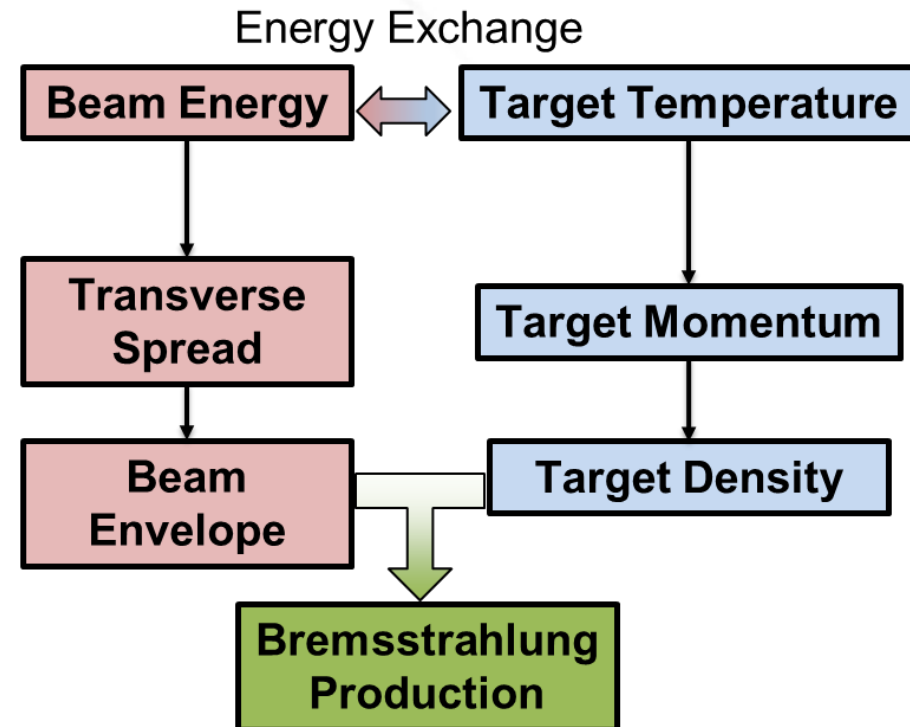
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Multi-pulse targets have intrinsic differences to single-pulse targets

- Target is evolving in time
- Input beam energy drives expansion
- X-ray output produced by co-incidence of beam and target
- Plume imaging provides insight into target state during evolution

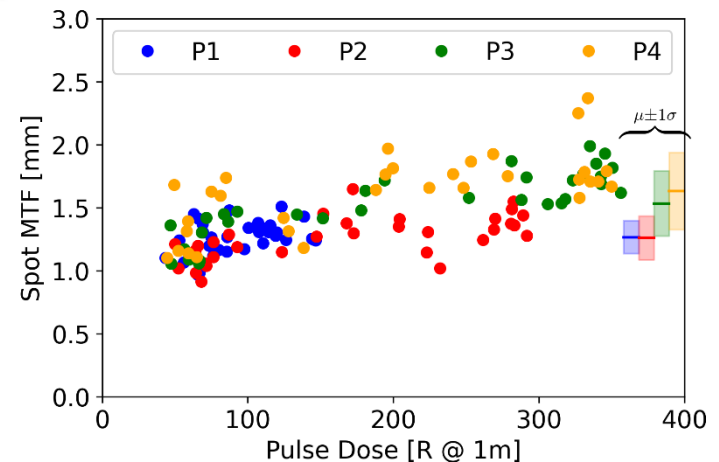


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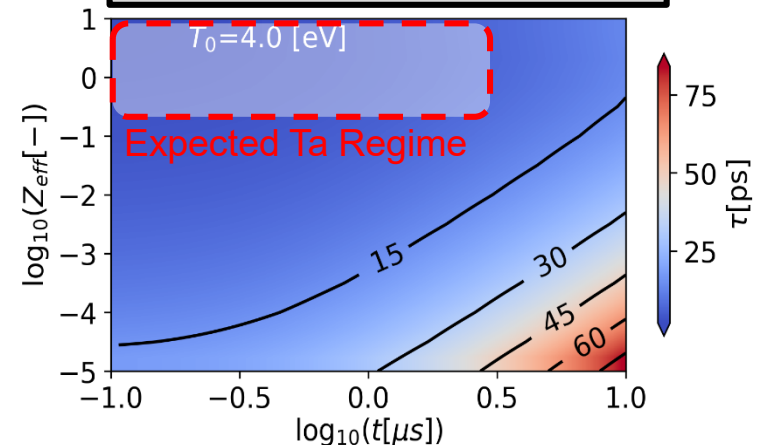
Beam evolution in target depends on target properties

- Historical database indicates late-time pulses are dissimilar to early times
- Beam-plasma interactions vary from macro- to micro-scale
 - Magnetic pinch effects evolve with magnetic diffusion time (density, temperature, degree of ionization)¹
 - Beam-plasma instability growth rate depends strongly on target density^{2,3}
- Understanding the target is integral to understanding our x-ray source

DARHT-2 Historical Performance



Oblique Instability Growth Rate



¹cf. Humphries "Charged Particle Beams" 2002.

²Bret, et al., Phys. Plasmas 17 (2010) 120501.

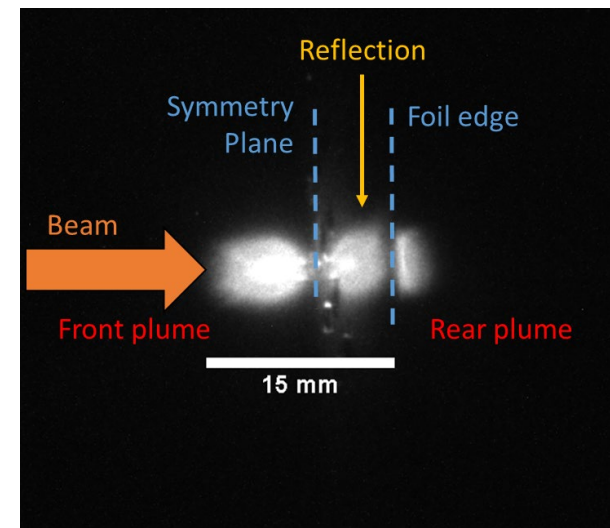
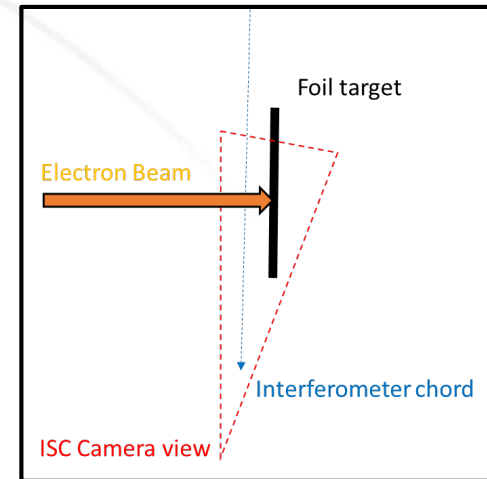
³Jaworski, et al., APS-DPP 2019, LA-UR-19-30587.

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Initial experiments conducted with single pulse driving expansion

- Camera positioned off-normal to avoid other optics
 - Angle is 5.5° and accounted for in analysis
 - Reflected image of plume is visible
- Timing of camera gate varied on identical tune
- Some camera artifacts present, but not used in analysis
- $250\mu\text{m}$ Mo foil target

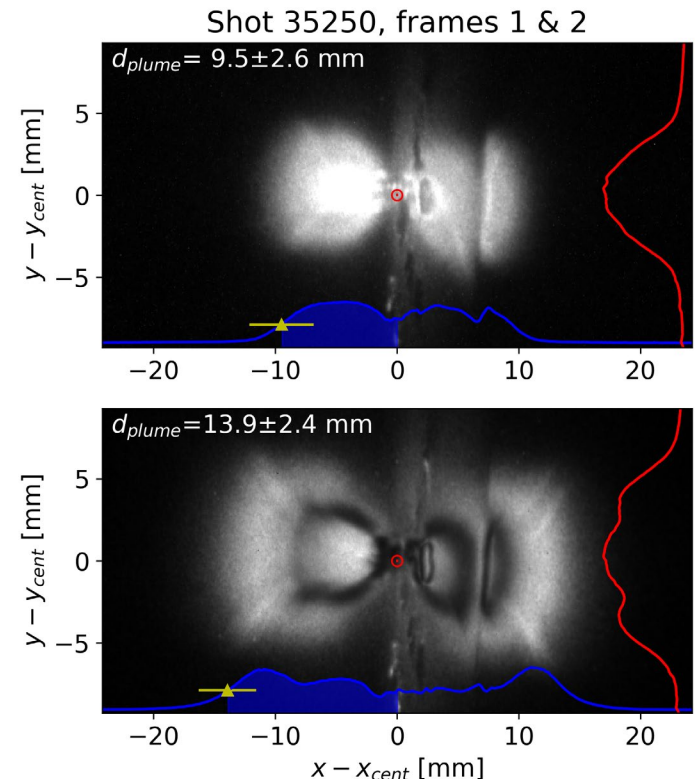


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Data reduction largely automated and includes uncertainty estimates

- Velocity extracted from each frame
 - Velocity = distance/time
 - Timing uncertainty based on earlier experiments scanning pulse (+/-10ns to t=0)
- Distance determined from center of plume (symmetry point)
 - Assumes photon intensity = material
 - Edges found at half-maximum intensity
- Uncertainty from multiple sources:
 - Intensity slope at half-max
 - Shifts of calculated center between frames
 - +/-1° viewing angle
 - Timing uncertainty



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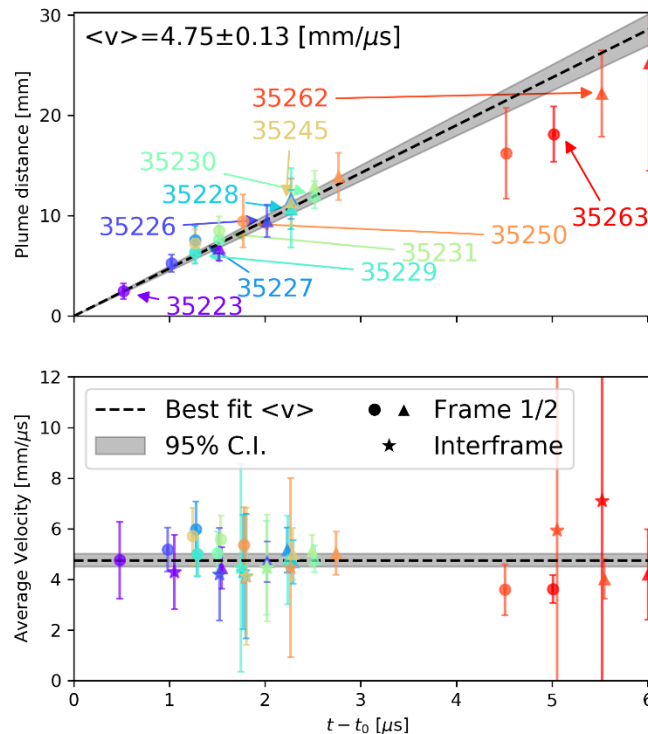
Plume expands at 4.75 mm/ μ s in this experiment

- From database select shots with:
 - S4 magnet = 100A, P1=40ns, P2=0ns
 - Frame 1 timing provides scan
- Constant velocity expansion model fit to data
 - Alternative models not explored
 - Potential deceleration late in time

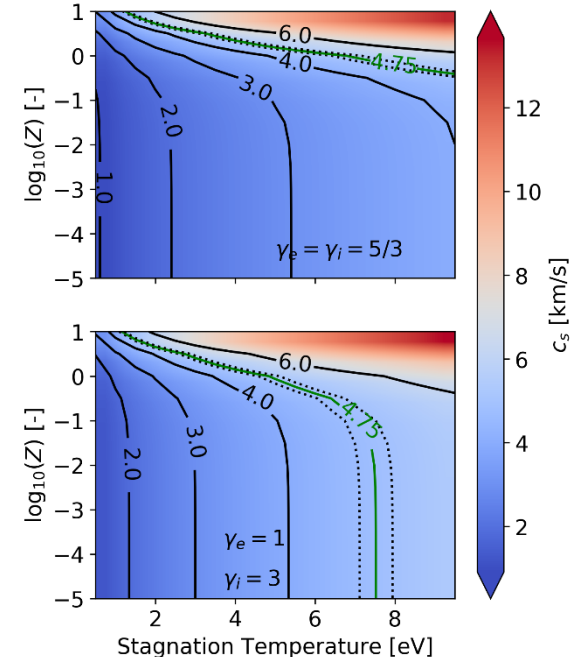
Ion Acoustic Speed

$$c_s = \left(\frac{\gamma_e Z_{eff} k T_e + \gamma_i k T_i}{M_i} \right)^{1/2}$$

Measured Velocities



Mo Sound Speed

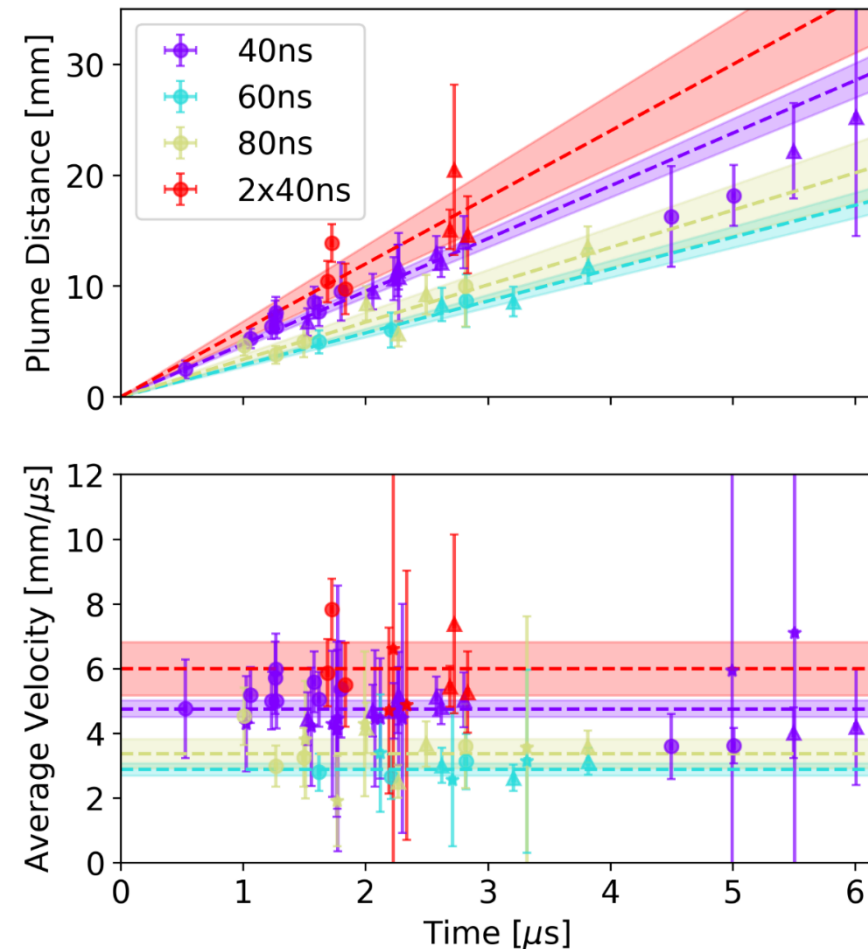


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Plume experiments examined four pulse formats

- 40, 60, 80, and 2x40ns experiments conducted
- Data indicate velocity relatively constant in all experiments
- Unexpected variations observed as pulse-length increases



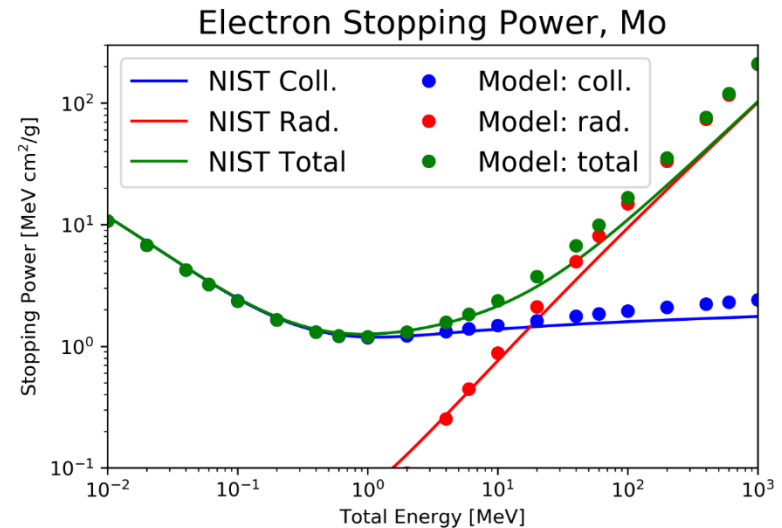
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E-beam parameters connect to thermodynamic model via energy density

- Spot radius, current, and pulse-length determine energy density
- Collisional stopping power only weakly dependent on energy
- Known parameters and material means radius can be extracted from known energy density

$$E_{dep} = \rho_{target} \frac{I_{beam}}{A_{beam}} \left. \frac{dE}{\rho dx} \right|_{coll.} \tau_{pulse}$$

$$r_{beam} = \left[\rho_{target} \frac{I_{beam}}{\pi E_{dep}} \left. \frac{dE}{\rho dx} \right|_{coll.} \tau_{pulse} \right]^{1/2}$$



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Mass, momentum, and energy balance used to evolve target material

- Mass balance: constant number of particles in the plume
- Momentum balance: change in velocity determined by pressure and mass in plume
- Energy balance: total energy ($e_{\text{int}} + e_{\text{kin}}$) is conserved except during heating
- Entropy constraints provide sanity-check on expansion
- **Real-gas thermodynamic model developed to describe pressure/temperature during expansion**

$$N_0 = n_0 h_0 \pi r_0^2 = n_0 V_0$$

$$n(t) = \frac{N_0}{V(t)}$$

$$F = m\vec{a} = \frac{d}{dt}(m\vec{v})$$

$$\Delta u = \frac{2P}{\rho d} \Delta t$$

$$E_{\text{total}} = E_{\text{kin.}} + E_{\text{int.}}$$

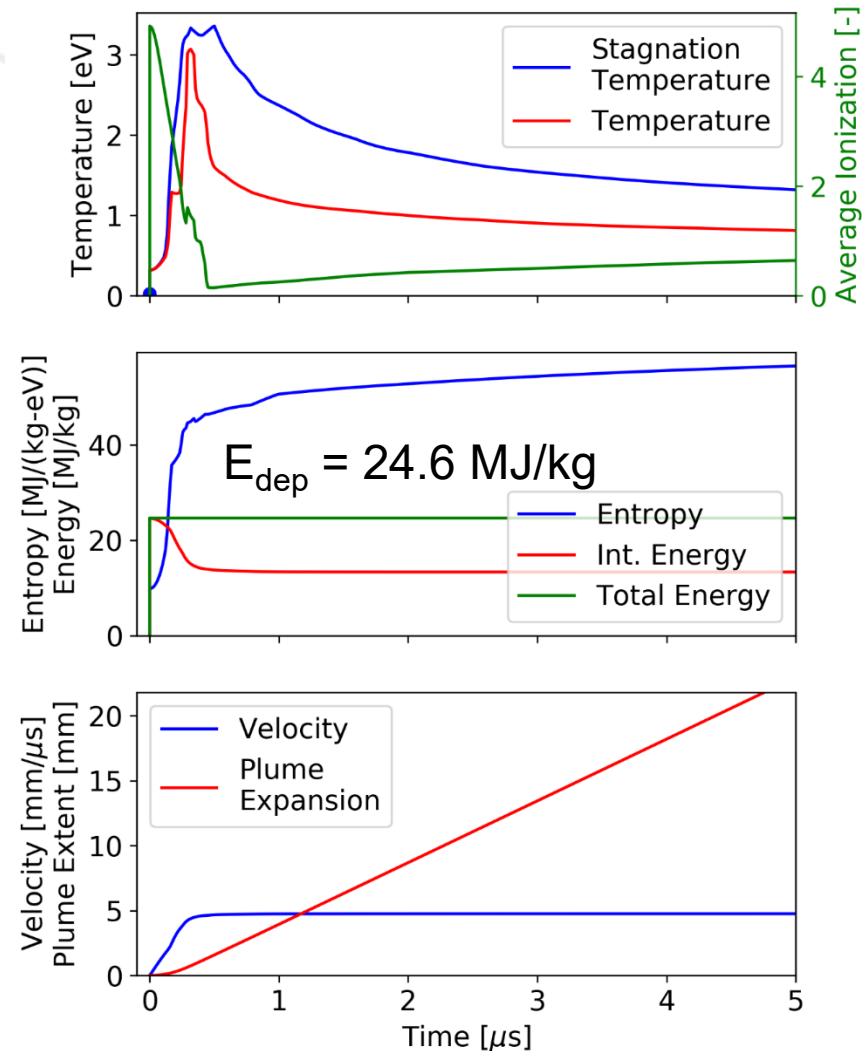
$$\Delta u \leq \frac{2}{\bar{u}} \left(\left. \frac{\partial s}{\partial \rho} \right|_e \right) \Delta \rho \left(\left. \frac{\partial s}{\partial e} \right|_\rho \right)^{-1}$$

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Thermodynamic model matches “terminal velocity” behavior

- Velocity saturates within $0.5\mu\text{s}$ (observed in experiment)
- Average ionization collapses while T increases (3-body recombination)¹
- Provides deposited energy consistent with these assumptions
- **IS IT CORRECT? Needs measurements!**



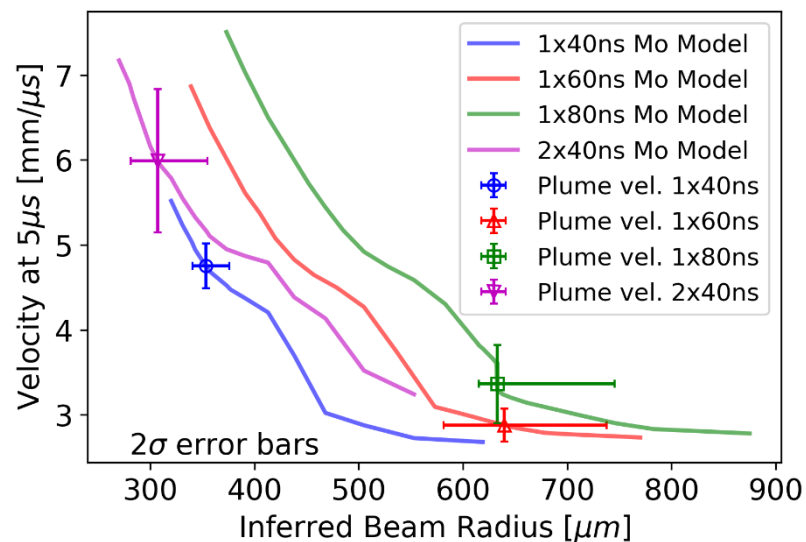
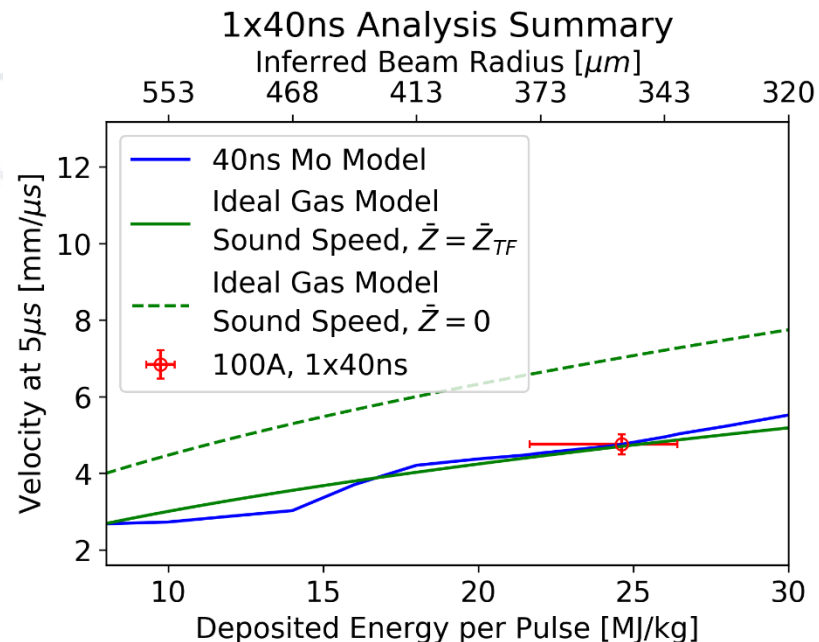
¹cf. Rumsby and Paul, Plasma Phys. 16 (1974) 247.

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Inferred spot sizes could indicate light-ion focusing effects

- Plume velocity uncertainty indicates range of E_{dep} in this model
- Axis 1 light-ion focusing effect measured with TRSS¹ showed reduction to $\frac{1}{2}$ integrated spot-size
- Compare to typical parameters:
 - FWHM $\sim 2 \times$ Radius
 - MTF $\sim 1.67 \times$ FWHM
 - $\sim 1.1\text{mm}$ MTF, $\sim 2\text{mm}$ MTF
- Validation of this type of model will lead to better understanding of beam-target interactions**



¹McCuistian, et al., Proc. EPAC08 (2008) 1206.

DARHT science stations provide significant utility for LIA operations and development

- OTR-based methods provide key parameters for the envelope equation
- OTR, solenoid scans, pepper-pot all provide ways to obtain \mathbf{R} , \mathbf{R}' , and ϵ_n
- Plume imaging is yielding insight into complex, late-time, target dynamics

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Thank you!

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